

MATRIX AND PROPERTIES

Matrix Type	Definition / Condition	Key Properties
Square Matrix	Same number of rows and columns	Can have determinant, trace, inverse, eigenvalues
Diagonal Matrix	All off-diagonal elements are zero	Diagonal elements may be non-zero
Scalar Matrix	Diagonal matrix with equal diagonal elements	Special case of diagonal matrix
Identity Matrix	Diagonal elements = 1; rest = 0	$A \cdot I = A, I \cdot A = A$
Zero Matrix	All elements are zero	$A + 0 = A, A \cdot 0 = 0$
Symmetric Matrix	$A^T = A$	Diagonalizable, real eigenvalues
Skew-Symmetric	$A^T = -A$	Diagonal elements = 0, eigenvalues are either 0 or purely imaginary
Hermitian Matrix	$A^\dagger = A$ (conjugate transpose = itself)	Complex symmetric analog, eigenvalues are real
Skew-Hermitian	$A^\dagger = -A$	Eigenvalues are purely imaginary or zero
Orthogonal Matrix	$A^T = A^{-1}$, i.e., $AA^T = I$	Preserves vector lengths, determinant = ± 1
Unitary Matrix	$A^\dagger = A^{-1}$, $AA^\dagger = I$	Complex analog of orthogonal matrix, preserves norm
Singular Matrix	Determinant = 0	Not invertible
Non-singular Matrix	Determinant $\neq 0$	Invertible
Triangular Matrix	All entries above or below diagonal are zero	Used in LU decomposition; determinant is product of diagonal elements
Normal Matrix	$AA^\dagger = A^\dagger A$	Diagonalizable using unitary matrices

PROBLEMS ON HERMITIAN, SKEW-HERMITIAN, SYMMETRIC, AND SKEW-SYMMETRIC MATRICES

No.	Problem	Solution
1	Determine whether the matrix $A = \begin{bmatrix} 2 & i \\ -i & 3 \end{bmatrix}$ is Hermitian.	Take conjugate transpose: $A^\dagger = \begin{bmatrix} 2 & i \\ -i & 3 \end{bmatrix}^\dagger = \begin{bmatrix} 2 & -i \\ i & 3 \end{bmatrix}$. This equals A , so Hermitian .
2	Is $A = \begin{bmatrix} 0 & 2i \\ -2i & 0 \end{bmatrix}$ a Skew-Hermitian matrix?	$A^\dagger = \begin{bmatrix} 0 & 2i \\ -2i & 0 \end{bmatrix}^\dagger = \begin{bmatrix} 0 & 2i \\ -2i & 0 \end{bmatrix}^* = \begin{bmatrix} 0 & -2i \\ 2i & 0 \end{bmatrix}$. Thus, $A^\dagger = -A \Rightarrow$ Skew-Hermitian .
3	Determine if $A = \begin{bmatrix} 1 & 2 \\ 2 & 3 \end{bmatrix}$ is symmetric.	$A^T = \begin{bmatrix} 1 & 2 \\ 2 & 3 \end{bmatrix} = A \Rightarrow$ Symmetric .
4	Find if $A = \begin{bmatrix} 0 & 5 \\ -5 & 0 \end{bmatrix}$ is skew-symmetric.	$A^T = \begin{bmatrix} 0 & -5 \\ 5 & 0 \end{bmatrix} = -A \Rightarrow$ Skew-symmetric .
5	Can a real skew-symmetric matrix have non-zero diagonal elements?	No, in skew-symmetric matrices, $a_{ii} = -a_{ii} \Rightarrow a_{ii} = 0$.
6	Show that the matrix $A = \begin{bmatrix} 0 & 2+i \\ -2-i & 0 \end{bmatrix}$ is skew-Hermitian.	$A^\dagger = \begin{bmatrix} 0 & -2+i \\ 2-i & 0 \end{bmatrix} = -A \Rightarrow$ Skew-Hermitian .
7	Show that every square matrix can be written as the sum of a Hermitian and a skew-Hermitian matrix.	For any matrix A , write: $A = \frac{1}{2}(A + A^\dagger) + \frac{1}{2}(A - A^\dagger)$. First part is Hermitian, second part is skew-Hermitian.
8	Check if the matrix $A = \begin{bmatrix} 1+i & 2 \\ 2 & 1-i \end{bmatrix}$ is Hermitian.	$A^\dagger = \begin{bmatrix} 1-i & 2 \\ 2 & 1+i \end{bmatrix} = A \Rightarrow$ Hermitian .
9	Find the Hermitian and skew-Hermitian parts of $A = \begin{bmatrix} 2 & 3+i \\ 3-i & 4 \end{bmatrix}$	Hermitian part $= \frac{1}{2}(A + A^\dagger) = \begin{bmatrix} 2 & 3 \\ 3 & 4 \end{bmatrix}$ Skew-Hermitian part $= \frac{1}{2}(A - A^\dagger) = \begin{bmatrix} 0 & i \\ -i & 0 \end{bmatrix}$
10	Determine if the matrix $A = \begin{bmatrix} 1 & 2i \\ -2i & 1 \end{bmatrix}$ is Hermitian.	$A^\dagger = \begin{bmatrix} 1 & 2i \\ -2i & 1 \end{bmatrix}^* = \begin{bmatrix} 1 & -2i \\ 2i & 1 \end{bmatrix}^T = \begin{bmatrix} 1 & 2i \\ -2i & 1 \end{bmatrix} = A$, hence Hermitian .

EIGENVALUES AND EIGENVECTORS

S.No.	Property	Explanation	Remarks/Examples
1	Definition	If $A\vec{x} = \lambda\vec{x}$, then λ is an eigenvalue and $\vec{x} \neq 0$ is an eigenvector.	$A \in \mathbb{R}^{n \times n}$, $\vec{x} \in \mathbb{R}^n$
2	Characteristic Equation	Eigenvalues are solutions of $\det(A - \lambda I) = 0$.	A polynomial equation of degree n .
3	Number of Eigenvalues	An $n \times n$ matrix has n eigenvalues (including complex and repeated).	Can be real or complex.
4	Linearly Independent Eigenvectors	Eigenvectors corresponding to distinct eigenvalues are linearly independent.	Important for diagonalization.
5	Eigenvalues of Diagonal Matrix	Are the diagonal entries themselves.	E.g. $D = \begin{bmatrix} 2 & 0 \\ 0 & 3 \end{bmatrix} \Rightarrow \lambda = 2, 3$
6	Eigenvalues of Triangular Matrix	Also lie on the diagonal.	Works for both upper and lower triangular matrices.
7	Trace	Sum of eigenvalues = Trace of matrix.	$\text{tr}(A) = \sum \lambda_i$
8	Determinant	Product of eigenvalues = Determinant of the matrix.	$\det(A) = \prod \lambda_i$
9	Eigenvalues of A^{-1}	If λ is an eigenvalue of A , then $\frac{1}{\lambda}$ is an eigenvalue of A^{-1} .	Provided A is invertible.
10	Eigenvalues of A^T	Same as those of A .	A and A^T have identical characteristic polynomials.
11	Eigenvalues of scalar multiple kA	Are $k\lambda_i$, if λ_i are eigenvalues of A .	Example: $A = \begin{bmatrix} 1 & 0 \\ 0 & 2 \end{bmatrix}$, $2A \Rightarrow \lambda = 2, 4$
12	Eigenvectors remain same under scalar multiplication	If $A\vec{x} = \lambda\vec{x}$, then $(kA)\vec{x} = k\lambda\vec{x}$.	Eigenvector is unchanged.
13	Eigenvalues of A^n	Are λ^n , if λ is an eigenvalue of A .	For powers of matrices.
14	Orthogonality (Symmetric Matrices)	Eigenvectors of real symmetric matrices are orthogonal.	Important in diagonalization.
15	Complex Eigenvalues	If A has real entries, then complex eigenvalues occur in conjugate pairs.	E.g. $\lambda = a + bi \Rightarrow \bar{\lambda} = a - bi$ also an eigenvalue.

BESSEL, LEGENDRE, HERMITE, AND LAGUERRE POLYNOMIALS

Property	Bessel $J_n(x)$	Legendre $P_n(x)$	Hermite $H_n(x)$	Laguerre $L_n(x)$
Type	Special function (non-polynomial)	Orthogonal polynomial	Orthogonal polynomial	Orthogonal polynomial
Order	Integer or real n	Integer n	Integer n	Integer n
Domain	$x \in \mathbb{R}, \mathbb{C}$	$x \in [-1,1]$	$x \in \mathbb{R}$	$x \in [0, \infty)$
Weight Function	$w(x) = x$	$w(x) = 1$	$w(x) = e^{-x^2}$	$w(x) = e^{-x}$
Orthogonality Relation	$\int_0^1 x J_n(\alpha_{n,m}x) J_n(\alpha_{n,k}x) dx = 0$ for $m \neq k$	$\int_{-1}^1 P_n(x) P_m(x) dx = 0$ for $n \neq m$	$\int_{-\infty}^{\infty} e^{-x^2} H_n(x) H_m(x) dx = 0$	$\int_0^{\infty} e^{-x} L_n(x) L_m(x) dx = 0$
Differential Equation	$x^2 y'' + xy' + (x^2 - n^2)y = 0$	$(1 - x^2)y'' - 2xy' + n(n + 1)y = 0$	$y'' - 2xy' + 2ny = 0$	$xy'' + (1 - x)y' + ny = 0$
Generating Function	$e^{\frac{x}{2}(t - \frac{1}{t})} = \sum_{n=-\infty}^{\infty} J_n(x)t^n$	$\frac{1}{\sqrt{1 - 2xt + t^2}} = \sum_{n=0}^{\infty} P_n(x)t^n$	$e^{2xt - t^2} = \sum_{n=0}^{\infty} \frac{H_n(x)}{n!} t^n$	$\frac{e^{-xt/(1-t)}}{1-t} = \sum_{n=0}^{\infty} L_n(x)t^n$
Recurrence Relation	$J_{n-1}(x) + J_{n+1}(x) = \frac{2n}{x} J_n(x)$	$(n + 1)P_{n+1} = (2n + 1)xP_n - nP_{n-1}$	$H_{n+1} = 2xH_n - 2nH_{n-1}$	$L_{n+1} = \frac{(2n + 1 - x)L_n - nL_{n-1}}{n + 1}$
Rodrigues' Formula	–	$P_n(x) = \frac{1}{2^n n!} \frac{d^n}{dx^n} [(x^2 - 1)^n]$	$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} [e^{-x^2}]$	$L_n(x) = \frac{e^x}{n!} \frac{d^n}{dx^n} [e^{-x} x^n]$
Applications	Cylindrical systems, EM waves	Spherical systems, potential theory	Quantum oscillator, probability	Hydrogen atom, quantum systems
Zeros	Infinite, real and positive	Real in $(-1,1)$	Real, symmetric about 0	Real, in $(0, \infty)$
Normalization	$J_0(0) = 1$	$P_n(1) = 1$	Leading term = 2^n	$L_n(0) = 1$

Property	Bessel $J_n(x)$	Legendre $P_n(x)$	Hermite $H_n(x)$	Laguerre $L_n(x)$
Terms	$J_0(x) = 1 - \frac{x^2}{4} + \frac{x^4}{64} - \frac{x^6}{2304} + \dots$ $J_1(x) = \frac{x}{2} - \frac{x^3}{16} + \frac{x^5}{384} - \dots$ $J_2(x) = \frac{x^2}{8} - \frac{x^4}{96} + \frac{x^6}{3072} - \dots$ $J_3(x) = \frac{x^3}{48} - \frac{x^5}{768} + \dots$ $J_4(x) = \frac{x^4}{384} - \frac{x^6}{9216} + \dots$ $J_5(x) = \frac{x^5}{3840} - \dots$	$P_0(x) = 1$ $P_1(x) = x$ $P_2(x) = \frac{1}{2}(3x^2 - 1)$ $P_3(x) = \frac{1}{2}(5x^3 - 3x)$ $P_4(x) = \frac{1}{8}(35x^4 - 30x^2 + 3)$ $P_5(x) = \frac{1}{8}(63x^5 - 70x^3 + 15x)$	$H_0(x) = 1$ $H_1(x) = 2x$ $H_2(x) = 4x^2 - 2$ $H_3(x) = 8x^3 - 12x$ $H_4(x) = 16x^4 - 48x^2 + 12$ $H_5(x) = 32x^5 - 160x^3 + 120x$	$L_0(x) = 1$ $L_1(x) = -x + 1$ $L_2(x) = \frac{1}{2}(x^2 - 4x + 2)$ $L_3(x) = \frac{1}{6}(-x^3 + 9x^2 - 18x + 6)$ $L_4(x) = \frac{1}{24}(x^4 - 16x^3 + 72x^2 - 96x + 24)$ $L_5(x) = \frac{1}{120}(-x^5 + 25x^4 - 200x^3 + 600x^2 - 600x + 120)$

Binomial, Poisson, and Normal Distributions

This guide compares the **Binomial**, **Poisson**, and **Normal** distributions in terms of their introduction, formulas, statistical moments, modes, characteristic functions, cumulants, cumulative functions, and applications.

Introduction

- **Binomial Distribution:** Describes the number of successes in a fixed number of independent Bernoulli trials (yes/no experiments) with the same probability of success.
- **Poisson Distribution:** Models the number of times an event occurs in a fixed interval of time or space, assuming events happen independently and at a constant mean rate.
- **Normal Distribution:** Represents a continuous probability distribution, known for its bell-shaped curve; it is used for distributions of sums (or averages) of many small, independent effects.

Distribution Formulas

Property	Binomial	Poisson	Normal
PDF/PMF	$P(X = k)$ $= \binom{n}{k} p^k (1 - p)^{n-k}$	$P(X = k)$ $= \frac{\lambda^k e^{-\lambda}}{k!}$	$f(x)$ $= \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{(x-\mu)^2}{2\sigma^2}}$
Parameters	n : trials, p : success prob.	λ : mean count	μ : mean, σ^2 : variance

Power Series Expansion

- **Binomial:** Based on binomial theorem, directly tied to discrete trials.
- **Poisson:** Derived as a limiting case of the binomial when $n \rightarrow \infty$, $p \rightarrow 0$, with $np = \lambda$ fixed.
- **Normal:** Taylor expansion in its density exponents; arises as the limiting case of binomial or Poisson by the Central Limit Theorem.

First Four Moments

Order	Binomial	Poisson	Normal
Mean	np	λ	μ
Variance	$np(1 - p)$	λ	σ^2
Skewness	$\frac{1 - 2p}{\sqrt{np(1 - p)}}$	$\frac{1}{\sqrt{\lambda}}$	0
Kurtosis (excess)	$\frac{1 - 6p(1 - p)}{np(1 - p)}$	$\frac{1}{\lambda}$	0

Mode(s)

- **Binomial:** $[(n + 1)p]$ (Mode is not always unique for integer values.)
- **Poisson:** $[\lambda]$
- **Normal:** The mode is at μ (also the mean and median).

Moment Generating Function (MGF)

Property	Formula
Binomial	$M(t) = (1 - p + pe^t)^n$
Poisson	$M(t) = \exp(\lambda(e^t - 1))$
Normal	$M(t) = \exp(\mu t + \frac{1}{2}\sigma^2 t^2)$

Cumulants

- **Binomial:** r th cumulant is $n \frac{d^r}{dt^r} \ln(1 - p + pe^t)|_{t=0}$

- **Poisson:** All cumulants are equal to λ .
- **Normal:** 1st is μ , 2nd is σ^2 , higher orders are zero.

Cumulative Distribution Function (CDF)

- **Binomial:** $F(k) = \sum_{i=0}^k \binom{n}{i} p^i (1 - p)^{n-i}$
- **Poisson:** $F(k) = \sum_{i=0}^k \frac{\lambda^i e^{-\lambda}}{i!}$
- **Normal:** $F(x) = \frac{1}{2} \left[1 + \operatorname{erf} \left(\frac{x - \mu}{\sigma\sqrt{2}} \right) \right]$ (No closed-form except using the error function.)

Typical Applications

Distribution	Typical Application Examples
Binomial	Quality control (defective products), coin flips, surveys
Poisson	Arrival of calls at a call center, radioactive decay, rare event modeling
Normal	Heights, test scores, measurement errors in science

Summary Table

Feature	Binomial	Poisson	Normal
Type	Discrete	Discrete	Continuous
Range	$0, 1, \dots, n$	$0, 1, \dots, \infty$	$(-\infty, \infty)$
Skewness	Varies with p	Decreases as λ increases	0 (symmetric)
Limiting Relation	-	Limit of Binomial	Limit of Binomial/Poisson as $n \rightarrow \infty$

MAXWELL'S EQUATIONS

Maxwell's Equations (in differential form)

Maxwell's equations unify electricity, magnetism, and optics into one electromagnetic theory. There are **four equations**, named after **James Clerk Maxwell**, based on experimental laws:

S.No.	Equation	Name	Empirical Basis
1.	$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$	Gauss's Law for Electricity	Coulomb's Law
2.	$\nabla \cdot \vec{B} = 0$	Gauss's Law for Magnetism	No magnetic monopoles
3.	$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$	Faraday's Law of Induction	Faraday's experiments
4.	$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$	Ampère–Maxwell Law	Ampère's Law + Maxwell's Displacement Current

◆ 1. Gauss's Law for Electricity

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

- **Meaning:** Electric field lines diverge from electric charges.
- ρ : charge density (C/m³)
- \vec{E} : Electric field (V/m)
- ϵ_0 : Permittivity of free space (8.854×10^{-12} F/m)

Units: LHS: $\nabla \cdot \vec{E} \Rightarrow \text{V/m}^2$ RHS: $\rho/\epsilon_0 \Rightarrow \text{C}/(\text{m}^3)/\text{F/m} = \text{V/m}^2$

◆ 2. Gauss's Law for Magnetism

$$\nabla \cdot \vec{B} = 0$$

- **Meaning:** No isolated magnetic monopoles; magnetic field lines are always **closed loops**.
- \vec{B} : Magnetic flux density (Tesla)

Units: LHS: $\nabla \cdot \vec{B} \Rightarrow \text{T/m} = \text{kg}/(\text{s}^2 \cdot \text{A})$ RHS: 0 (dimensionless) — implies conservation.

◆ 3. Faraday's Law of Induction

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

- **Meaning:** Time-varying magnetic fields induce **electric fields** (EMF).
- Basis: Faraday's experimental observations

Units: LHS: $\nabla \times \vec{E} \Rightarrow \text{V/m}^2$ RHS: $-\partial \vec{B} / \partial t \Rightarrow \text{T/s} = \text{V/m}^2$

◆ 4. Ampère–Maxwell Law

$$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

- **Meaning:** Magnetic fields are produced by **currents and changing electric fields**
- First term: Ampère's Law
- Second term: **Displacement current** added by Maxwell

Units:

- LHS: $\nabla \times \vec{B} \Rightarrow \text{T/m} = \text{kg}/(\text{s}^2 \cdot \text{A})$
- RHS:
 - $\mu_0 \vec{J}$: (H/m) · A/m² = T/m
 - $\mu_0 \epsilon_0 \partial \vec{E} / \partial t = \text{T/m}$

📖 Summary Table with Units

Equation	Vector Form	Physical Meaning	Units
1. Gauss's Law (E)	$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$	Electric flux divergence equals charge density	V/m ²
2. Gauss's Law (B)	$\nabla \cdot \vec{B} = 0$	No magnetic monopoles	T/m
3. Faraday's Law	$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$	Time-varying B induces E	V/m ²
4. Ampère–Maxwell Law	$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$	Currents and time-varying E induce B	T/m

Mechanics of Particles

Mechanics of particles deals with the motion and interaction of particles (objects with mass but negligible size). It includes the study of **force, momentum, energy**, and their **conservation laws**.

Conservation of Linear Momentum

Principle:

In the absence of external force, the **total linear momentum** of a system remains **constant**.

$$\vec{p}_{\text{initial}} = \vec{p}_{\text{final}} \quad \text{if } \vec{F}_{\text{ext}} = 0$$

Example:

A **0.05 kg bullet** is fired from a **1 kg gun** with velocity $\mathbf{v} = 100 \text{ m/s}$. Find the recoil velocity of the gun.

Solution:

Initial momentum = 0

Final momentum: $0 = m_{\text{gun}} \cdot v_{\text{gun}} + m_{\text{bullet}} \cdot v_{\text{bullet}}$

$$0 = 1 \cdot v_{\text{gun}} + 0.05 \cdot 100 \Rightarrow v_{\text{gun}} = -5 \text{ m/s}$$

Gun recoils at 5 m/s in the opposite direction.

Conservation of Angular Momentum

Principle:

If no external torque acts on a system, its **angular momentum** remains constant. $L = \vec{r} \times \vec{p}, \quad \vec{\tau}_{\text{ext}} = 0 \Rightarrow \frac{dL}{dt} = 0$

Example:

A skater has **moment of inertia** $I = 4 \text{ kg}\cdot\text{m}^2$ and **angular velocity** $\omega = 2 \text{ rad/s}$. When arms are pulled in, $I = 2 \text{ kg}\cdot\text{m}^2$. What is the new angular velocity?

Solution:

$$L_1 = L_2 \Rightarrow I_1 \omega_1 = I_2 \omega_2 \Rightarrow 4 \cdot 2 = 2 \cdot \omega_2 \Rightarrow \omega_2 = 4 \text{ rad/s}$$

New angular velocity = 4 rad/s.

Conservation of Energy

Principle:

In a conservative system (no friction or loss), **mechanical energy** (kinetic + potential) is conserved. $E = KE + PE = \text{constant}$

Example:

A **2 kg** object falls from **10 m**. What is its speed just before hitting the ground?

Solution:

$$\text{Initial PE} = mgh = 2 \cdot 9.8 \cdot 10 = 196 \text{ J} \quad \text{Final KE} = \frac{1}{2}mv^2 = 196 \Rightarrow v =$$

14 m/s

Speed = 14 m/s.

QUICK REVISION POINTS

Mechanics of Particles and Systems of Particles

- **Particle:** Idealized object with mass but no size.
- **System of particles:** A collection of interacting particles.
- **Newton's laws** apply individually to each particle.
- **Centre of mass:** Average position of the mass in a system.
- **Linear momentum (P)** of system = total mass \times velocity of center of mass.
- **Angular momentum (L)** of system = conserved if net external torque is zero.
- **Total energy** = Kinetic + Potential (conserved in isolated systems).

□ Constraints and Generalized Coordinates

- **Constraint:** Restriction on motion (e.g., a bead on a wire).
- **Generalized coordinates (q_1, q_2, \dots, q_n):** Minimal set of independent coordinates to describe a system.
- **Degrees of freedom:** Number of independent coordinates.
- **Holonomic system:** Constraints can be expressed as equations (e.g., $f(q, t) = 0$).
- **Non-holonomic:** Inequalities or differential constraints (e.g., rolling without slipping).
- **Scleronomic:** Constraints independent of time.
- **Rheonomic:** Constraints explicitly time-dependent.

💡 Conservation Laws

- **Conservation of Energy:** Total mechanical energy is conserved if the force is conservative.
- **Conservation of Linear Momentum:** Conserved if **no external force** acts on the system.
- **Conservation of Angular Momentum:** Conserved if **no external torque** acts.
- **Conservative forces:** Derivable from potential energy $\vec{F} = -\nabla V$.
- **Non-conservative forces:** Path-dependent (e.g., friction).

□ Lagrangian Formalism

- **Lagrangian:** $L = T - V$, where T is kinetic and V is potential energy.
- **Lagrange's Equations:**

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0$$

- **D'Alembert's Principle:** Introduces virtual work to derive equations of motion.

🔍 Applications:

- **Simple Pendulum:** One degree of freedom (angle θ).
- **Atwood's Machine:** Two-mass pulley system.
- **Harmonic Oscillator:** Restoring force proportional to displacement.
- **Electrical LC Circuit:** Analogy: Charge \leftrightarrow Position, Current \leftrightarrow Velocity.

□ Hamiltonian Formalism

- **Hamiltonian:** $H = \sum_i p_i \dot{q}_i - L$, often equals **total energy**.
- **Hamilton's Equations:**

$$\dot{q}_i = \frac{\partial H}{\partial p_i}, \quad \dot{p}_i = -\frac{\partial H}{\partial q_i}$$

- **Cyclic (ignorable) coordinates:** If $\partial L / \partial q_i = 0$, then p_i is conserved.

- **Variational Principle (Hamilton's Principle):** Motion follows a path that minimizes the **action**:

$$\delta \int_{t_1}^{t_2} L dt = 0$$

- **Principle of Least Action:** Foundation of all classical mechanics.
- **Canonical Transformations:** Preserve Hamiltonian structure of equations.
- **Liouville's Theorem:** Phase space volume is conserved in time.

🌀 Rigid Body Dynamics

- **Rigid Body:** Distance between particles remains constant.
- **Moment of Inertia Tensor:** Describes how mass is distributed with respect to rotational axes.
- **Euler's Angles (θ, ϕ, ψ):** Specify orientation of body in space.

📖 Euler's Equations of Motion

- For a rigid body rotating about a point:

$$\begin{aligned} I_1 \dot{\omega}_1 + (I_3 - I_2) \omega_2 \omega_3 &= N_1 \\ I_2 \dot{\omega}_2 + (I_1 - I_3) \omega_3 \omega_1 &= N_2 \\ I_3 \dot{\omega}_3 + (I_2 - I_1) \omega_1 \omega_2 &= N_3 \end{aligned}$$

- Derived in body-fixed frame (rotating axes).

☐ Symmetrical Top under Gravity

- Symmetric: $I_1 = I_2 \neq I_3$
- Fixed point at one end, center of mass offset.
- Uses **Euler angles**.
- Motions:
 - **Spin:** About symmetry axis.
 - **Precession:** Around vertical axis.
 - **Nutation:** Wobbling of axis.

- Lagrangian:

$$L = \frac{1}{2} I_1 (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + \frac{1}{2} I_3 (\dot{\psi} + \dot{\phi} \cos \theta)^2 - mgl \cos \theta$$